

# Progress and Challenges in Specifying Geomagnetic Activity for the Electrical Power Grid

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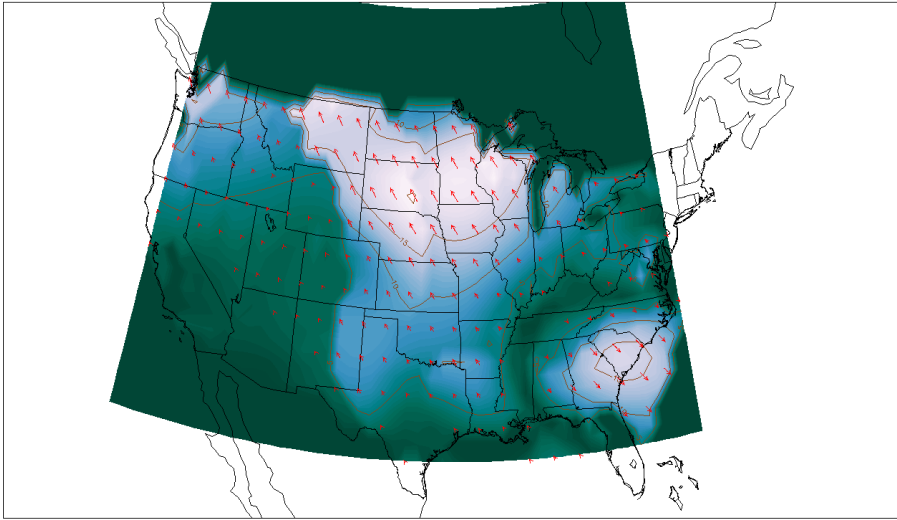
## Outline

- Long-term objective
- Drivers versus response functions
- Normalizing out the response
- The way forward

# Long-term Objective

Geoelectric field map

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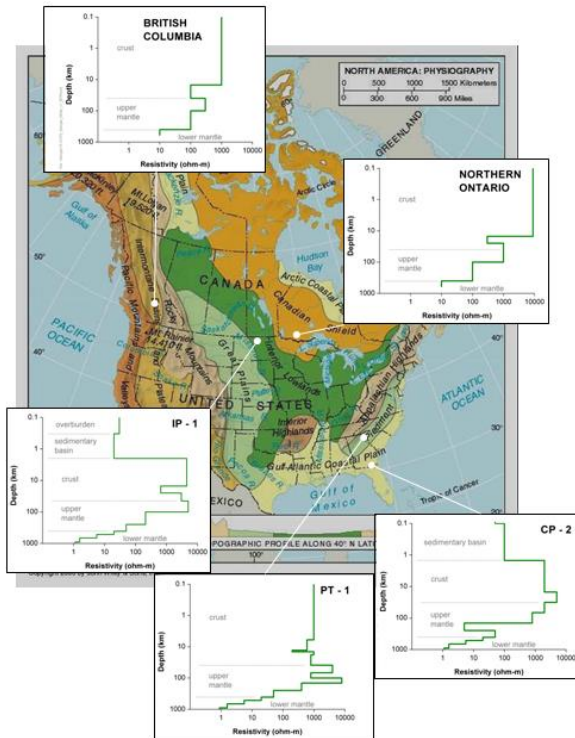
**Operational product to calculate the local geoelectric field (V/km)**

- **Use real-time magnetometer data as the input**
- **Interpolate geomagnetic field variations on a grid between observatories**
- **Calculate Electric field 'locally' using appropriate conductivity models**
- **User applies the electric field to a model for the power grid to calculate geomagnetically induced currents**
- **User assesses system stability and transformer vulnerability**

## Two key components

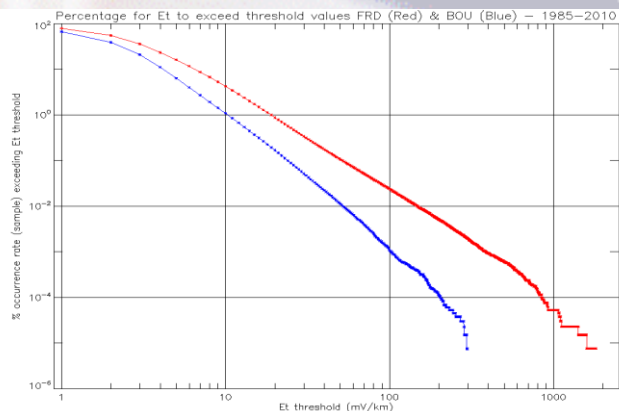
- **The External Driver (Space Weather)**
  - Time varying currents in the ionosphere and magnetosphere driven by interaction with disturbed solar wind
- **The Geological Conductivity Structure**
  - Naturally induced currents below Earth's surface
  - Significantly modifies impact of Space Wx driver

# Initial Efforts - 1D Conductivity Models



Credit: Fernberg, Gannon and Bedrosian

- 1D Conductivity profiles provided by USGS for ~20 different physiographic regions
- Preliminary 1-minute values for Ex & Ey have been calculated for 1985-2012
- Seven USGS observatories
- Overall histograms
- Storm Profiles
- Climatology conditioned on existing measures (Kp)
- Validation work I.P.



# What about that Conductivity Model?

- **Good News:**
  - Once you figure it out, it won't change
- **Bad News:**
  - It is very complex, inhomogeneous, highly structured, and not always well known
- **Advice of our partners at USGS:**
  - 1D models for the physiographic regions are probably not sufficiently accurate

## **Normalized E-field: Uniform Half Space**

- **Uniform half-space solution can be used to ‘normalize’ the geology**
- **Succinctly characterizes the external driving component of the E-field**
- **As specific conductivity models improve, the normalized E-field can be modified fairly easily to incorporate this information**
- **Users may be able to use their own GIC data and system models to empirically determine the ‘Earth Transfer function’ without knowing the details of the conductivity**

# Uniform half-space

- Conductivity  $\sigma_c$ , resistivity  $\rho_c = 1/\sigma_c$
- Time dependence  $e^{i\omega t}$
- Plane wave solution in conducting medium:

$$E_x = E_0 e^{-k_c z}, \text{ where } k_c = \sqrt{i\omega\mu\sigma_c}$$

- Faraday's law:  $\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} = -\mu \frac{\partial \mathbf{H}}{\partial t}$

- Polarized Plane wave

$$-k_c E_x = -i\omega\mu H_y$$

- Hence the electric field is related to the magnetic field by

$$E_x = \frac{i\omega\mu}{k_c} H_y$$

- But using the form for  $k_c$ , above we find

$$E_x = \sqrt{i\omega\mu\rho_c} H_y$$

## Uniform Half Space (cont)

- $\sqrt{i\omega\mu\rho_c}$  - 'Earth Transfer Function' for a uniform half-space
- $\omega^{1/2}$  shows the primary dependence of the E-field on  $\frac{\partial B}{\partial t}$ 
  - Relatively higher frequencies in B get an extra boost from the earth transfer function
  - Greatest interest is in the 0.001 to 0.01 Hz range
- A normalized value of  $\rho_c = 1$  ohm-m can be used to compare the input driver uniformly over the entire continent
- Local calculations have to correct for the effect of the local conductivity – will show the scaling for some simple cases

# Earth Transfer Function: Multi-layer Model

- n layers ( $i=0,1,2,\dots,n-1$ ), last layer semi-infinite
- Conductivities  $\sigma_i$ , depths  $h_i$
- $k_i = \sqrt{i\omega\mu\sigma_i}$
- For each layer:

$$E_x = A_i e^{-k_i z} + B_i e^{k_i z}$$

$$H_y = \frac{k_i}{i\mu\omega} (A_i e^{-k_i z} - B_i e^{k_i z})$$

- Define impedance  $Z_i$  as ratio of  $E_x/H_y$  at  $z=h_i$
- Derive a recurrence relationship between  $Z_i$  and  $Z_{i+1}$
- General form:

$$Z_i = \frac{i\omega\mu}{k_{i+1}} \left( \frac{1 - \alpha_{i+1}}{1 + \alpha_{i+1}} \right)$$

## Multi-layer solution

- Recurrence calculation repeated layer by layer to get the surface impedance:

$$Z_S = \frac{i\omega\mu}{k_0} \left( \frac{1-\alpha_0}{1+\alpha_0} \right) = \sqrt{i\omega\mu\rho_0} \left( \frac{1-\alpha_0}{1+\alpha_0} \right)$$

# Comparison

- Uniform half-space

$$E_x = \sqrt{i\omega\mu\rho_c} H_y$$

- Multi-layer solution:

$$E_x = \sqrt{i\omega\mu\rho_0} \left( \frac{1-\alpha_0}{1+\alpha_0} \right) H_y$$

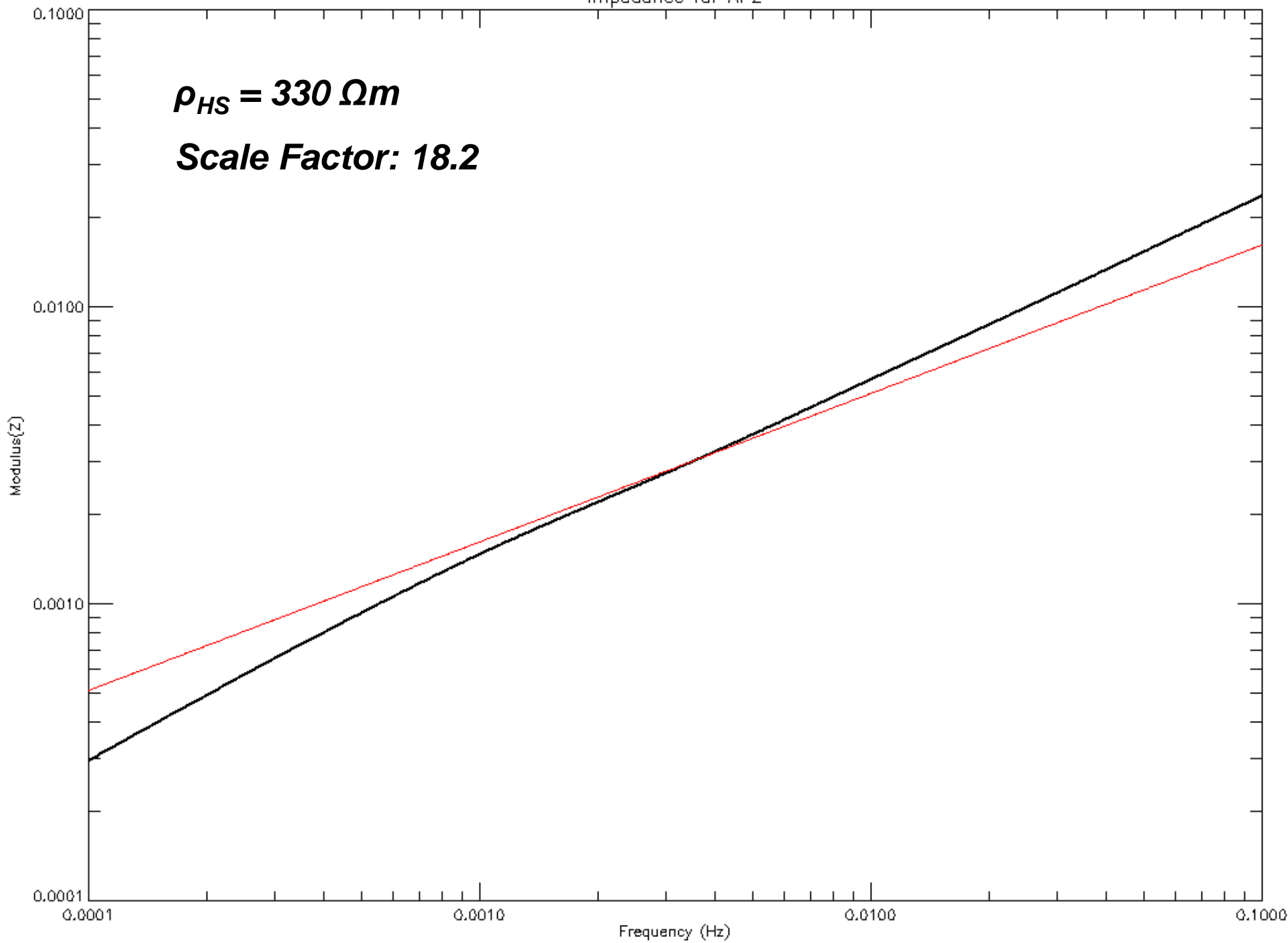
- The relationship between the two:

$$E_x(layer) = \sqrt{\frac{\rho_0}{\rho_c}} \left( \frac{1-\alpha_0}{1+\alpha_0} \right) E_x(halfspace)$$

- A comparison of the impedances for uniform half space versus multi-layer solutions can help sort out the relative role of the driver versus the geology

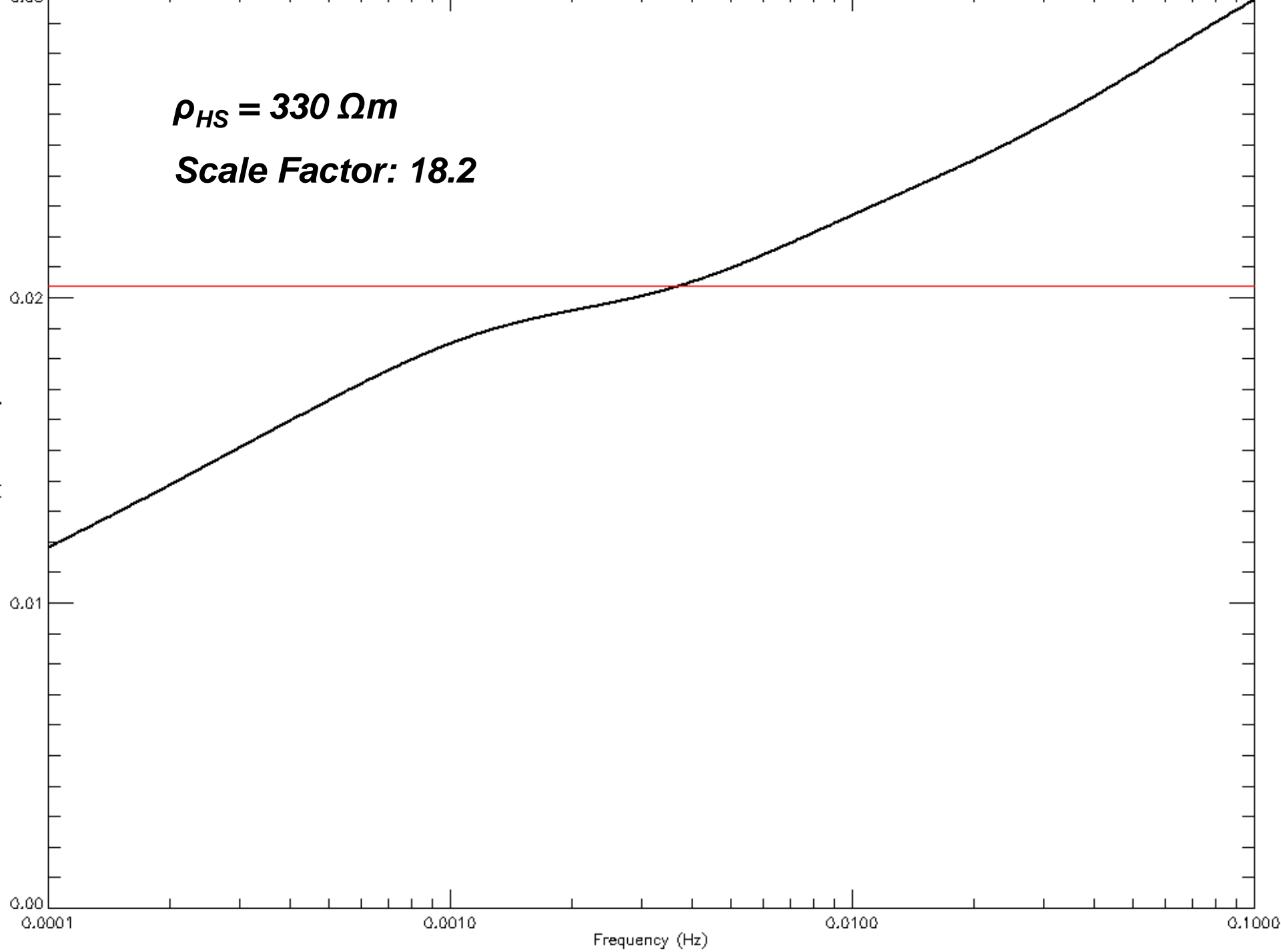
$$\rho_{HS} = 330 \Omega m$$

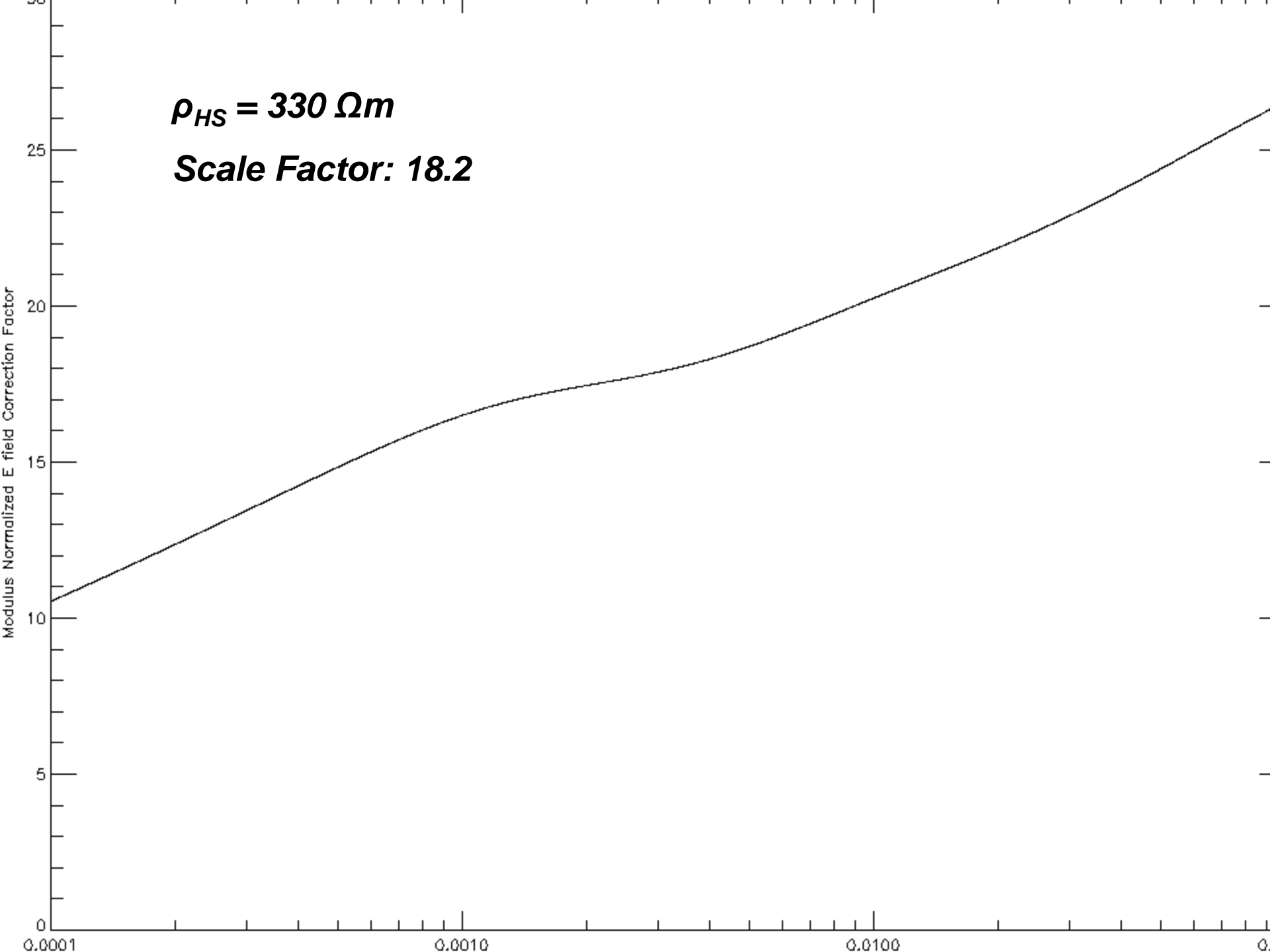
**Scale Factor: 18.2**



**$\rho_{HS} = 330 \Omega m$**

**Scale Factor: 18.2**





# Uniform Half-Space Fitting

Model	$\rho_{HS} (\Omega m)$	Scaling	Model	$\rho_{HS} (\Omega m)$	Scaling
PT1	972	31	IP1	179	13
CP2min	676	26	PB1	123	11
CP2	671	26	BC	98	10
OTT	548	23	AK1A	87	9
CP1	491	22	AK1B	87	9
CL1	466	22	SL1	86	9
SU1	463	21	CS1	44	7
IP3	463	22	AP1	40	6
CP2max	377	19	PB2	37	6
AP2	330	18	CO1	36	6
NE1	314	18	IP4	26	5
IP2	257	16	BR1	22	5

- Ordered by resistivity, highest to lowest
- Scaling factor gives approximate correction to earth-transfer function
- The conductivity profile can change the answer by a factor of 6!

*$\rho_{HS}$  is the resistivity of a uniform half-space that gives the best fit to the multi-layer model over 0.001 to 0.01 Hz*

# The way forward...

- **Working jointly with NOAA, USGS and NRCAN for new operational product development**
- **Validation of the 1D E-field values is in process (NASA/CCMC leading the comparisons)**
  - Looking at accuracy & consistency of techniques
  - Looking for accuracy of the numbers where possible
- **Priority to derive  $B(r,t)$  on a spatial grid using existing network with flexibility to expand**
  - There is a collaborative effort being led by EPRI and DOE to deploy additional magnetometers to improve the accuracy of  $B(r,t)$  interpolation

# The way forward...

- We are considering providing a 'generic' normalized  $E(r,t)$  on a spatial grid using a uniform half space conductivity
- There would be an option to provide the corrected calculation from a catalogue of 1D profiles
- User requirement for 0.001 to 0.010 Hz is driving the need to obtain higher time-resolution magnetometer data
  - 1 minute data only goes to 0.008 Hz
  - Analysis needed to assess the error this causes

# Summary

- **Electric grid users need the geoelectric field**
- **The Space Weather part of the calculation is succinctly derived from geomagnetic field data using uniform half-space conductivity model**
- **The conductivity part of the calculation is difficult but significant – options are:**
  - **Use a 1D profile (your own or the USGS physiographic regions)**
  - **Work with the scientific community to improve the conductivity specification in a particular region**
  - **Use GIC data and system models to derive empirical estimates for the earth transfer function**

# Multi-layer solution

- At the bottom of layer  $i+1$ :

$$Z_{i+1} = \frac{E_x(h_{i+1})}{H_y(h_{i+1})} = \frac{i\omega\mu}{k_{i+1}} \left( \frac{A_{i+1}e^{-k_{i+1}h_{i+1}} + B_{i+1}e^{k_{i+1}h_{i+1}}}{A_{i+1}e^{-k_{i+1}h_{i+1}} - B_{i+1}e^{k_{i+1}h_{i+1}}} \right)$$

- For notation simplification define:

$$C_{i+1} = \frac{Z_{i+1}}{i\omega\mu}$$

- Then the ratio of B to A for layer  $i+1$  is found to be:

$$R_{i+1} = \frac{B_{i+1}}{A_{i+1}} = - \left( \frac{1 - k_{i+1}C_{i+1}}{1 + k_{i+1}C_{i+1}} \right) e^{-2k_{i+1}h_{i+1}}$$

- At the top of layer  $i+1$ :

$$Z_i = \frac{E_x(h_i)}{H_y(h_i)} = \frac{i\omega\mu}{k_{i+1}} \left( \frac{A_{i+1}e^{-k_{i+1}h_i} + B_{i+1}e^{k_{i+1}h_i}}{A_{i+1}e^{-k_{i+1}h_i} - B_{i+1}e^{k_{i+1}h_i}} \right)$$

# Multi-layer solution

- Using the ratio of  $B_{i+1}$  to  $A_{i+1}$  we get :

$$Z_i = \frac{i\omega\mu}{k_{i+1}} \left( \frac{(1 + k_{i+1}C_{i+1}) - (1 - k_{i+1}C_{i+1})e^{-2k_{i+1}(h_{i+1}-h_i)}}{(1 + k_{i+1}C_{i+1}) + (1 - k_{i+1}C_{i+1})e^{-2k_{i+1}(h_{i+1}-h_i)}} \right)$$

- We can define parameter  $\alpha_{i+1}$  (using  $d_i$  as the thickness)

$$\alpha_{i+1} = \left( \frac{1 - k_{i+1}C_{i+1}}{1 + k_{i+1}C_{i+1}} \right) e^{-2k_{i+1}d_{i+1}} \text{ (note that } d_{N-1} \rightarrow \infty, \alpha_{N-1} = 0)$$

- Resulting in the recurrence relation:

$$Z_i = \frac{i\omega\mu}{k_{i+1}} \left( \frac{1 - \alpha_{i+1}}{1 + \alpha_{i+1}} \right)$$

## Multi-layer solution

- Top of semi-infinte layer (bottom of layer N-2):

$$Z_{N-2} = \frac{i\omega\mu}{k_{N-1}}, \text{ thus } C_{N-2} = \frac{Z_{N-2}}{i\omega\mu}$$

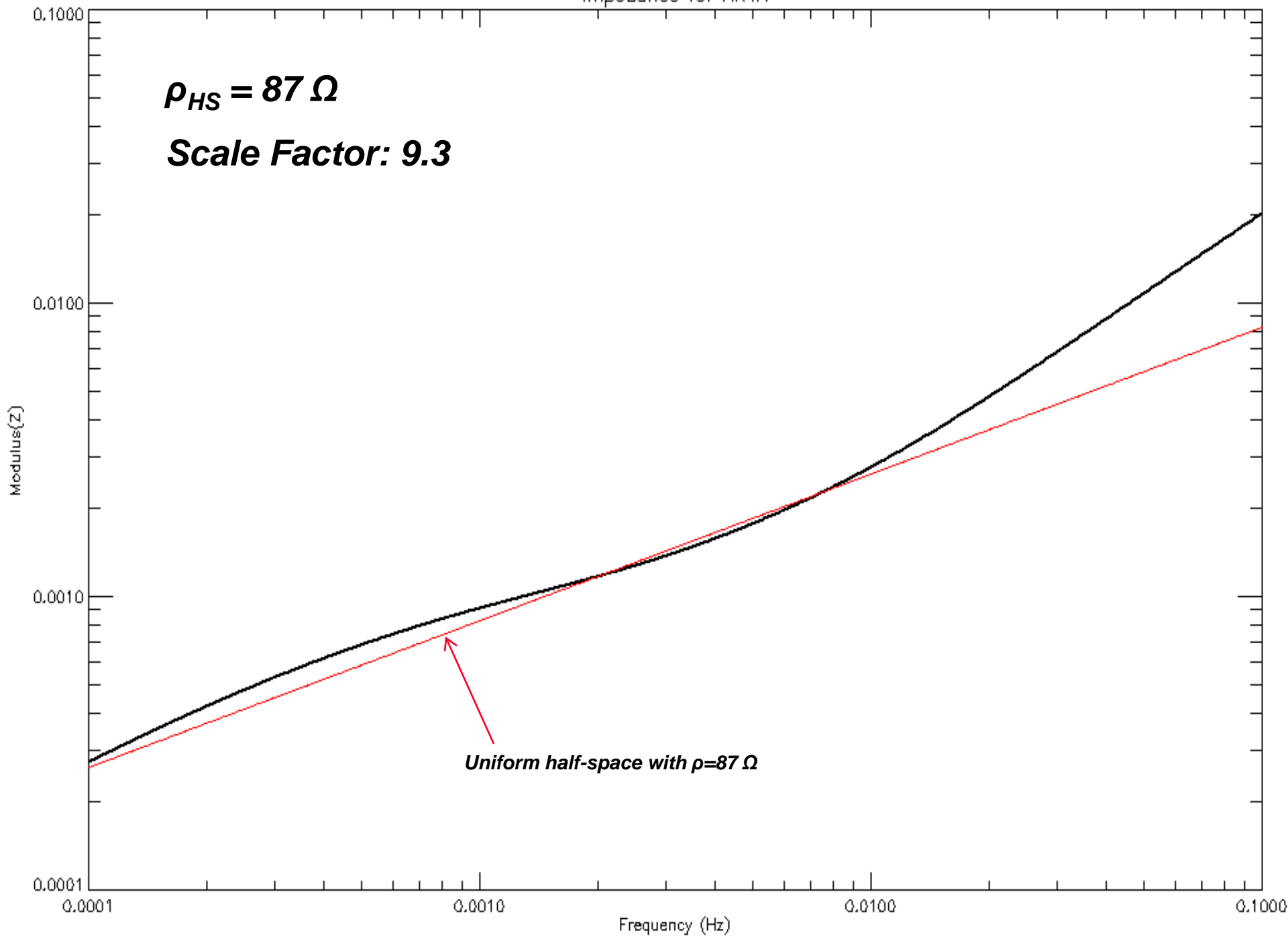
- Calculate  $\alpha_{N-2} = \left( \frac{1 - k_{N-2} C_{N-2}}{1 + k_{N-2} C_{N-2}} \right) e^{-2k_{N-2} d_{N-2}}$

- Find impedance at top of next layer (bottom of layer N-3):

$$Z_{N-3} = \frac{i\omega\mu}{k_{N-2}} \left( \frac{1 - \alpha_{N-2}}{1 + \alpha_{N-2}} \right)$$

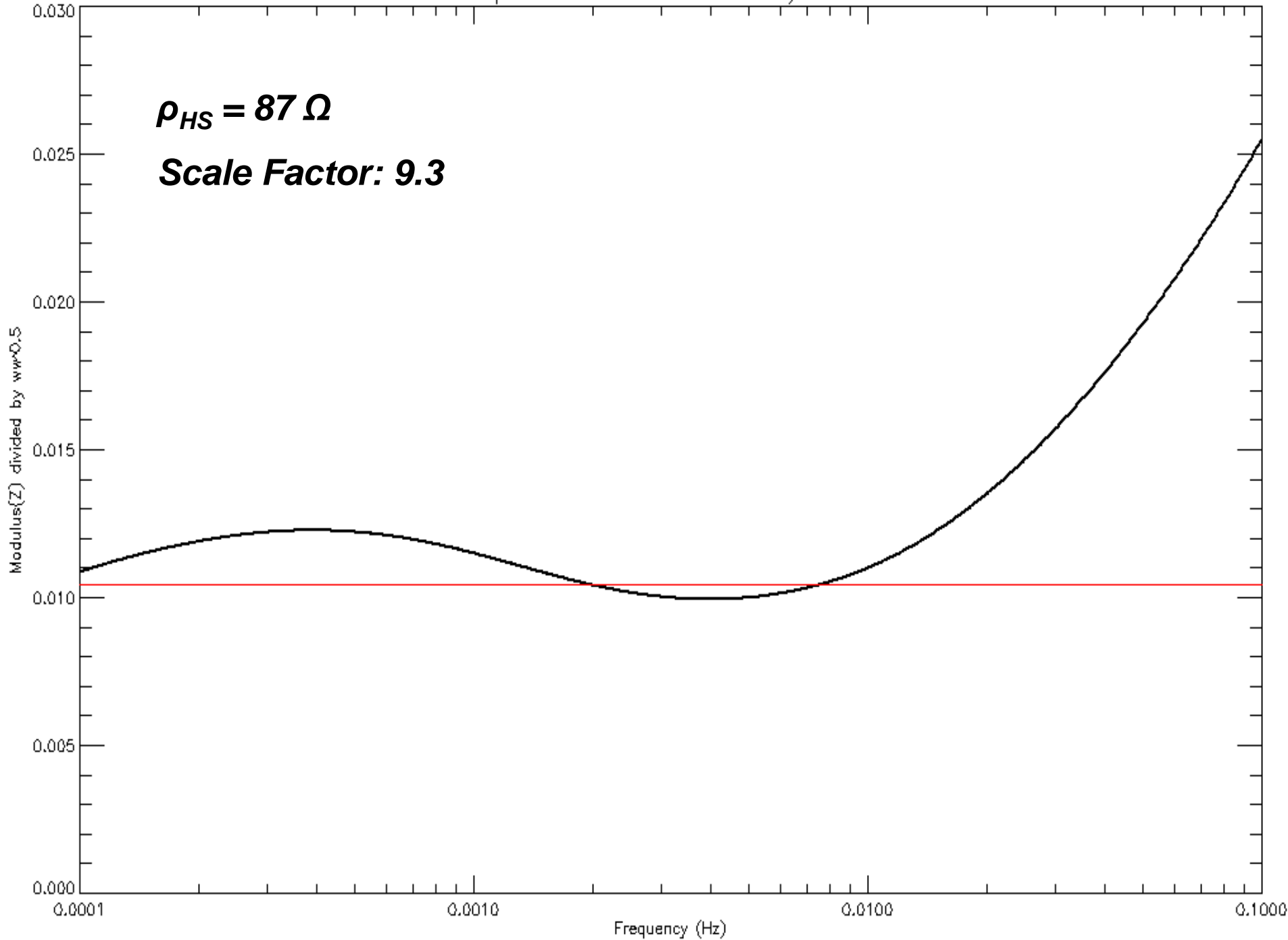
- etc, up to  $Z_0, \alpha_0$ . Finally we get the surface impedance:

$$Z_S = \frac{i\omega\mu}{k_0} \left( \frac{1 - \alpha_0}{1 + \alpha_0} \right) = \sqrt{i\omega\mu\rho_0} \left( \frac{1 - \alpha_0}{1 + \alpha_0} \right)$$



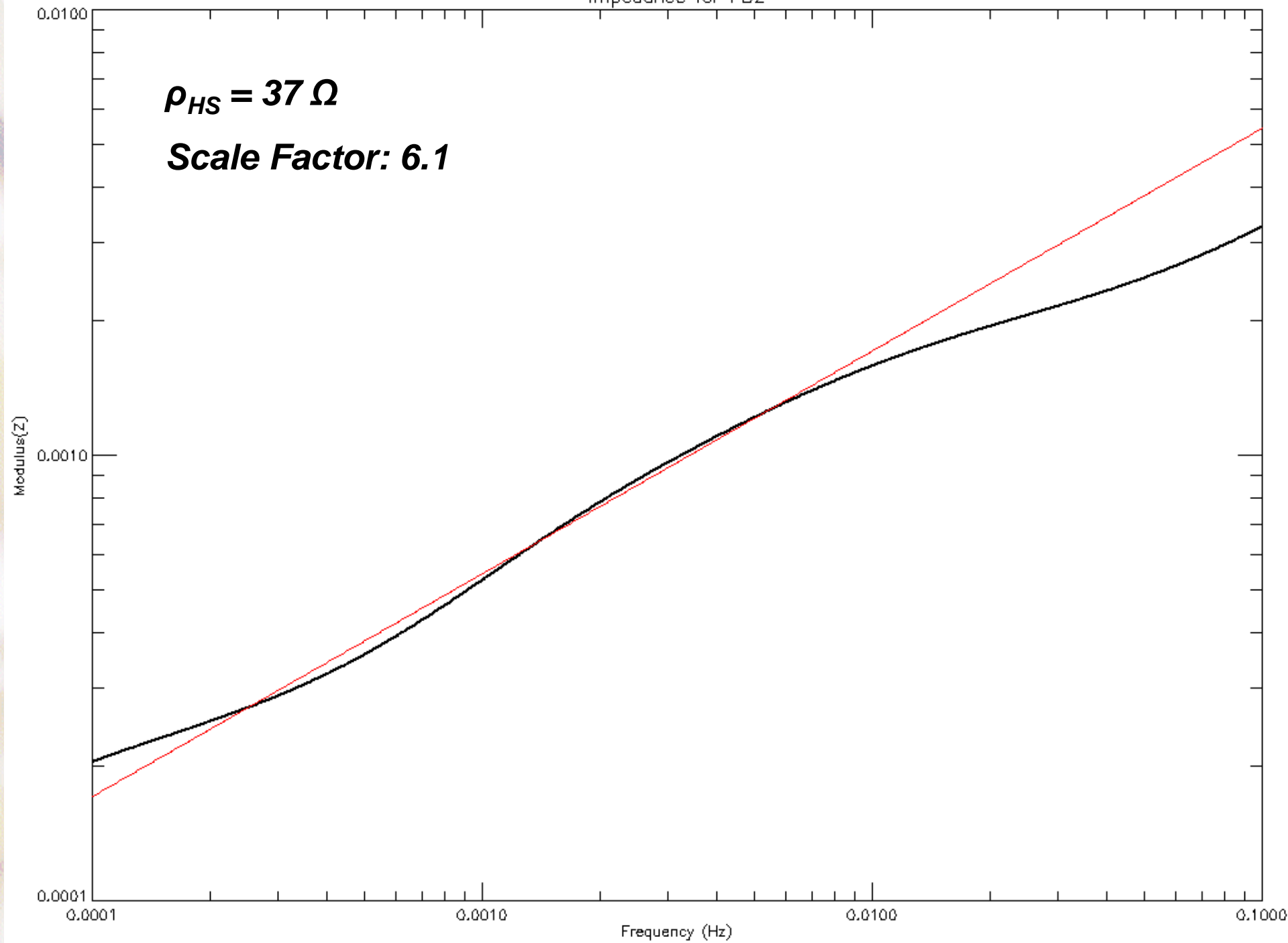
**$\rho_{HS} = 87 \Omega$**

**Scale Factor: 9.3**



$$\rho_{HS} = 37 \, \Omega$$

**Scale Factor: 6.1**



$$\rho_{HS} = 37 \, \Omega$$

**Scale Factor: 6.1**

